Variational Principles and Lagrange's Equations Lecture 3

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Hamilton's Variational Principle

Statement- Hamilton's variational principle for conservative system is stated as follows. The motion of system from time t_1 to t_2 is such that the line integral,

$$I = \int_{t_1}^{t_2} L dt \tag{1}$$

is extremum for the path of motion. where L = T - V is called the Lagrangian.

Let us consider a conservative system of particles. Employing the generalised co-ordinates, the integral can be written as;

$$\int_{t_1}^{t_2} [T(q_j, \dot{q}_j) - V(q_j)] dt$$
 (2)

Then according to Hamilton's variational principle,

$$\delta \int_{t_1}^{t_2} [T(q_j, \dot{q}_j) - V(q_j)] dt = 0$$
(3)

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$$\int_{t_1}^{t_2} \sum_{j} \left[\left(\frac{\partial T}{\partial q_j} \delta q_j + \frac{\partial T}{\partial \dot{q}_j} \delta \dot{q}_j \right) - \frac{\partial V}{\partial q_j} \right] dt = 0$$

$$\int_{t_1}^{t_2} \sum_{j} \left(\frac{\partial T}{\partial q_j} - \frac{\partial V}{\partial q_j} \right) \delta q_j dt + \int_{t_1}^{t_2} \sum_{j} \frac{\partial T}{\partial \dot{q}_j} \delta \dot{q}_j dt = 0$$
(4)

$$\int_{t_1}^{t_2} \sum_j \left(\frac{\partial T}{\partial q_j} - \frac{\partial V}{\partial q_j}\right) \delta q_j dt + \int_{t_1}^{t_2} \sum_j \frac{\partial T}{\partial \dot{q}_j} \frac{d}{dt} (\delta q_j) dt = 0$$
(6)

$$\int_{t_1}^{t_2} \sum_j \left(\frac{\partial T}{\partial q_j} - \frac{\partial V}{\partial q_j}\right) \delta q_j dt + \sum_j \left.\frac{\partial T}{\partial \dot{q}_j} \delta q_j\right|_{t_1}^{t_2} - \int_{t_1}^{t_2} \sum_j \left.\frac{d}{dt} \left(\frac{\partial T}{\partial \dot{q}_j}\right) \delta q_j dt = 0$$
(7)

In such variation $\delta q_j \Big|_{t_1}^{t_2}$, Hence

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$$\int_{t_1}^{t_2} \sum_{j} \left(\frac{\partial T}{\partial q_j} - \frac{\partial V}{\partial q_j}\right) \delta q_j dt - \int_{t_1}^{t_2} \sum_{j} \frac{d}{dt} \left(\frac{\partial T}{\partial \dot{q}_j}\right) \delta q_j dt = 0$$

$$\int_{t_1}^{t_2} \sum_{j} \left[\frac{\partial T}{\partial q_j} - \frac{\partial V}{\partial q_j} - \frac{d}{dt} \left(\frac{\partial T}{\partial \dot{q}_j}\right)\right] \delta q_j dt = 0$$
(8)

Since all δq_j are independent of each other the coefficient of δq_j should be zero.

$$\frac{\partial T}{\partial q_j} - \frac{\partial V}{\partial q_j} - \frac{d}{dt} \left(\frac{\partial T}{\partial \dot{q}_j} \right) = 0$$
(10)

$$\frac{d}{dt}\left(\frac{\partial T}{\partial \dot{q}_j}\right) - \frac{\partial}{\partial q_j}(T - V) = 0$$
(11)

For conservative system V is not the function of \dot{q}_j . Hence

$$\frac{d}{dt}\left(\frac{\partial}{\partial \dot{q}_j}(T-V)\right) - \frac{\partial}{\partial q_j}(T-V) = 0$$
(12)

$$\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{q}_j}\right) - \frac{\partial L}{\partial q_j} = 0 \tag{13}$$

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Non-holonomic system (Lagrange's method of undetermined multipliers)

The non-holonomic constraints can be put in the form,

$$\sum_{k} a_{lk} dq_k + a_{lt} dt = 0, l = 1, 2, 3...$$
(14)

The Lagrange's equations are then be obtained by inserting above equation in Hamilton's variational principle. This is the extension of Hamilton's variational principle for non-holonomic systems. This procedure is called Lagrange's method of undetermined multipliers. The virtual displacement in Hamilton's variational principle are taken at constant time and so above equation becomes,

$$\sum_{k} a_{lk} dq_k = 0 \tag{15}$$

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If there are m constants we have m equations in all l = 1, 2, 3...m. Multiply eq (15) by m constants $\lambda_i = \lambda_1, \lambda_2, \lambda_3,\lambda_m$ sum over l and integrate resulting equation from point 1 to 2.

$$\int_{1}^{2} \sum_{k} \sum_{l} \lambda_{l} a_{lk} \delta q_{k} dt = 0$$
(16)

Combine this with Hamilton's principle for conservative systems as below;

$$\delta \int_{1}^{2} L dt = \int_{1}^{2} dt \sum_{k} \left(\frac{\partial L}{\partial q_{k}} - \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_{k}} \right) \right) \delta q_{k} = 0$$
(17)
$$\int_{1}^{2} dt \sum_{k} \left(\frac{\partial L}{\partial q_{k}} - \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_{k}} \right) + \sum_{l} \lambda_{l} a_{lk} \right) \delta q_{k} = 0$$
(18)

All the $q'_k s$ are not independent but are connected by m equations (15). However the first (n-m) of these co-ordinates may be chosen independently the last m ones being fixed by eq (15). Let us choose λ_l 's to be such that

$$\frac{\partial L}{\partial q_k} - \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_k} \right) + \sum_l \lambda_l a_{lk} = 0$$
(19)

where k = n - m + 1, ..., n Having found λ_l 's from (19) we can write eq(18)

$$\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{q}_k}\right) - \frac{\partial L}{\partial q_k} = \sum_l \lambda_l a_{lk}$$
(20)

This is the set of Lagrange's equations for non-holonomic systems.

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